Influence of the CO₂ Antisolvent Effect on Ultrasound-Induced Polymer Scission Kinetics

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ABSTRACT: Ultrasound-induced polymer scission is a nonrandom process which alters the molecular weight distribution of polymers. However, transient cavitation, and consequently polymer scission, is not possible in concentrated polymer solutions due to the high liquid viscosity. The addition of an antisolvent can be used to circumvent this problem because the antisolvent decreases the gyration radius of polymer chains, which induces a reduction in liquid viscosity. To determine the influence of carbon dioxide (CO₂) as an antisolvent on the ultrasound-induced scission rate, ultrasonic scission experiments of poly(methyl methacrylate) have been performed in bulk methyl methacrylate (MMA) as well as in CO₂-expanded MMA. Modeling the experimental time-dependent molecular weight distributions (MWD) has revealed the scission kinetics at different polymer concentrations and CO₂ fractions. At low polymer concentrations, the scission rate is decreased upon an increased CO₂ content. This is a result of the higher vapor pressure of CO₂, which cushions the cavitation. However, at higher polymer concentrations, this effect is counteracted by the viscosity reduction induced by CO₂. As a consequence, the scission rate in CO₂-expanded MMA is higher as compared to bulk MMA for solutions with a high polymer concentration. The results show that ultrasound-induced scission in pressurized CO₂ can alter and control the MWD of polymers even in concentrated polymer solutions, whereas ultrasound-induced scission in bulk solutions is limited to relatively low polymer concentrations.

Introduction

The molecular weight distribution (MWD) is an important characteristic of polymers. In the polymer industry, a postprocessing step is often applied to alter the MWD, e.g., peroxide-induced degradation of polypropylene. In this process, fracture of the polymer chains occurs at random sites. An alternative method is ultrasound-induced polymer scission, which involves a much better controlled, nonrandom process. It can either be used as a postprocessing step or can be used during ultrasound-induced polymerizations.

Ultrasound is known to enhance chemical reactions, as well as mass transfer, at ambient temperatures and pressures. Most of these effects are caused by transient cavitations, i.e., the collapse of microscopic bubbles in a liquid. The chemical effects of cavitation result from the extreme conditions in the bubble (5000 K and 200 bar)⁴ and the high strain rates outside the bubble (10⁷ s⁻¹)⁵ generated upon implosion. Polymer scission arises from the high strain rates near an imploding cavity. The chain is fractured approximately in the middle of the chain, until a limiting molecular weight (M_{lim}) is reached.⁶ The M_{lim} is defined as the maximum chain length that does not break by ultrasound. A direct consequence of the nonrandom scission behavior is that no oligomers are formed by ultrasound-induced scission, in contrast to chemical degradation of polymers. Oligomers often have a negative influence on the polymer properties and are thus undesired in the final product. Additionally, an interesting application of ultrasound-induced polymer scission is the production of block copolymers, for which two different methods are possible. The first route

involves the dissolution of a homopolymer in a different monomer. Subsequently, ultrasonic scission of the polymer chains will generate the radicals that initiate the polymerization reaction with the monomer present. In the second method, two different homopolymers are dissolved in a nonreactive solvent. Ultrasound generates polymeric radicals, which will terminate by cross combination. The advantage of this method is that block copolymers can be produced from homopolymers of which the polymer—monomer systems are immiscible. For the system considered in this study, poly(methyl methacrylate) (PMMA) radicals react with a radical scavenger. In experiments without radical scavenger present, PMMA radicals would terminate by disproportionation.

Typically, polymer scission experiments are performed in the corresponding monomer as the solvent. Several studies have considered the solvent quality during scission experiments, 9,10 for which an increased degradation rate and a decreased $M_{\rm lim}$ have been observed when better solvents are used. In the literature, these effects have been ascribed to the configuration of the polymer chain, which becomes more open in better solvents. Additionally, some correlation with the χ parameter has been observed.

The viscosity is an important factor during ultrasound-induced scission. 11 A high viscosity hinders both the growth of a cavity, as well as the collapse of a cavity, and consequently, the polymer scission rate is reduced. At increasing polymer concentration, the scission process is less effective and eventually stops, as a high viscosity hinders the bubble implosion. This is a drawback for the development of a scission process based on ultrasound, as concentrated polymer systems are favored in industry. The addition of an antisolvent for the polymer can counteract the increase in viscosity at

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higher polymer concentrations. ¹² At relatively low antisolvent concentrations, the viscosity reduction is a result of the smaller gyration radius of the dissolved polymer molecules. It is still a one-phase system, due to which, the scission rate constant is expected to remain virtually unchanged. At higher CO₂ fractions, the polymer precipitates, and as a result, a second phase is formed. The precipitated polymer cannot be stretched by cavitation, and subsequently, it is no longer susceptible for ultrasound-induced scission. A constant scission rate is thus expected with increasing polymer concentration, as the liquid viscosity is unchanged by precipitation of the polymer.

High-pressure CO₂ is an interesting medium for ultrasound-induced polymer scission because it exhibits an antisolvent effect for most polymers, whereas most monomers have a high solubility in CO₂. ¹³ As a consequence, a CO₂-based process allows for relative easy separation of the polymer product from the reaction medium. Moreover, CO₂ is regarded as an environmentally friendly compound that is nontoxic, nonflammable, and naturally abundant. 14,15 However, up till now, ultrasound has rarely been studied at higher pressures because in most cases a high static pressure hampers the growth of cavities. Recently, we have shown that transient cavitation is possible in dense-phase fluids. 16 This provides a means for the development of scission processes based on ultrasound and high-pressure CO₂. These types of processes can contribute to a safer and cleaner way to produce polymers with the desired MWD. By adjusting the CO₂ fraction in the liquid, the amount of polymer that is dissolved can be controlled and, hence, the scission rate. Moreover, the CO₂ fraction influences the liquid viscosity and the vapor pressure, which determine the cavitation intensity.2 Therefore, the cavitational collapse, and consequently, the fracture rate and the M_{lim} are dependent on the amount of CO_2 in the liquid.

In this work, the influence of the CO₂ antisolvent effect on the ultrasound-induced polymer scission has been studied. Scission experiments have been performed at different CO₂ fractions and polymer concentrations. With the development of the experimental MWDs, as well as a general scission model based on population balances, the influence of CO₂ and the polymer concentration on the scission kinetics have been determined.

Ultrasound-Induced Scission. The breakage of polymer molecules due to ultrasound is a direct result of transient cavitation because no polymer fracture occurs when cavitation is suppressed. ¹⁷ Upon bubble collapse, the entire molecule moves along with the fluid. Due to the velocity profile near the cavities, friction between the polymer chain and the liquid occurs (Figure 1). This friction can result in stretching of the polymer coil and subsequent fracture of the chain. As described in a previous paper, the breaking of the polymer chain appears to be the rate-limiting step, not the stretching of the chain. 18 Since the drag force accumulates in the middle of the chain, scission primarily occurs at the center because the chance to break is the largest in the middle of the polymer chain. Of course, multiple breakage can occur during one scission experiment. However, it is not possible to remove a small fragment from a long polymer chain by ultrasound-induced scission because the forces acting on that bond are too small. The chain is fractured if the accumulated drag force in a chain fragment is higher than the bond strength. This force

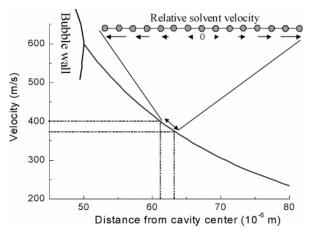


Figure 1. Schematic representation of the velocity profile near an imploding bubble and the resulting relative liquid velocity near a stretched polymer chain.

(eq 1) is dependent on the size of the monomer molecules (a and b represent the radius and the length of a monomeric unit in a polymer chain, ¹⁸ respectively), the shielding factor (S), the solvent viscosity (η), the number of monomer units (n), and the strain rate (d ϵ /dt). ¹⁹ The strain rate can be calculated from the distance of the polymer chain from the center of the cavity (r), the bubble wall velocity (dR/dt), and the radius of the bubble (R) (eq 2). ²⁰

$$F_{c} = \frac{6\pi}{8} \eta ab S \frac{\mathrm{d}\epsilon}{\mathrm{d}t} n^{2} \tag{1}$$

$$\frac{\mathrm{d}\epsilon}{\mathrm{d}t} = -2\frac{R^2 \frac{\mathrm{d}R}{\mathrm{d}t}}{r^3} \tag{2}$$

The two most important quantities describing an ultrasound-induced scission reaction are the $M_{
m lim}$ and the scission rate. The first is determined by the maximum strain rate in the liquid, which is obtained at the cavity interface at the moment of implosion. From the cavitation bubble interface toward the bulk liquid, the strain rate decreases rapidly. Polymer chains with a molecular weight slightly higher than $M_{\rm lim}$ can only break in the vicinity of the bubble interface, whereas higher-molecular-weight polymers can also break further away from the cavitation bubble. The scission rate is thus proportional to the volume around the cavity in which a polymer chain can stretch and break. Since the strain rate and the length of the polymer chain determine this volume, relatively low-molecular-weight polymers have a small volume in which they can break and, hence, have a low scission rate. The scission volume increases for polymer chains with a higher molecular weight and for faster imploding bubbles. The kinetics of ultrasonic scission can thus be directly ascribed to the implosion velocity of the cavitation bubbles and the length of the polymer that is fractured. The implosion velocity is dependent on the liquid properties, the content of the bubble, the acoustic pressure, and the static pressure, respectively. 21 Therefore, it is important to know the phase composition, density, and viscosity of the liquid.

Assuming first-order scission kinetics, the scission rate can be determined as a function of molecular weight

and liquid properties. The kinetic scheme of scission by ultrasound reads as follows²²

Scission:
$$P_n \xrightarrow{k_d(n,m)} P_m^{\bullet} + P_{n-m}^{\bullet}$$

Termination by

recombination: $P_m^{\bullet} + P_{n-m}^{\bullet} \xrightarrow{k_{tc}} P_n$

Termination by

disproportionation:
$$P_m^{\bullet} + P_{n-m}^{\bullet} \xrightarrow{k_{td}} P_m + P_{n-m}$$

If a radical scavenger is present, the termination reaction occurs instantaneously only between a macroradical fragment and a scavenger molecule. Thus, the scheme reduces to

$$P_n \xrightarrow{k_{\mathrm{d}}(n,m)} P_m + P_{n-m}$$

where $k_d(n,m)$ is an effective scission rate constant. The scission rate constant $k_{\rm d}$ is chosen as a function of the original chain length, n, and fragment length, m, to allow for the nonrandom character of the scission process. According to the overall scission reaction, the population balance then becomes

$$\frac{dP_n}{dt} = \sum_{m=n+1}^{\infty} k_{d}(n,m)P_m - P_n \sum_{m=1}^{n} k_{d}(n,m)$$
 (3)

This balance equation is solved with the Galerkin finite element (FEM) package PREDICI.²³ By manipulating the shape of $k_d(n,m)$, which is anticipated to be Gaussian with a width (σ), it is possible to obtain a best fit between the MWD predicted by eq 4 and the measured MWD. The overall dissociation constant k_d consists of a Gaussian distribution function and a molecular-weight-dependent scission function (k_{MW}) :¹⁸

$$k_{\mathrm{d}}(n,m) = \frac{k_{\mathrm{MW}}(n)}{\sqrt{2\pi}\sigma(n)} \mathrm{e}^{\left(-\frac{1}{2\sigma(n)^2}\left(m - \frac{n}{2}\right)^2\right)} \tag{4}$$

As the force in the middle of the chain has a squared dependence in the number of monomeric units (eq 1), an almost quadratic function is used to describe the dissociation rate as a function of the molecular weight (eq 5). This equation is only valid for polymer chains with a molecular weight higher or equal to $M_{
m lim}$. When the dissociation constant k_{MW} is zero, the corresponding $M_{\rm lim}$ at these conditions is obtained. This molecular weight is determined by the constants A and B. The scission rate is defined by all three constants in eq 5:18

$$k_{\text{MW}}(n) = A + Bn^{2e^{-n/C}} \tag{5}$$

Experimental Section

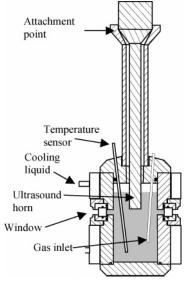
Materials. The polymer solutions used for the scission experiments with ultrasound consisted of MMA (Merck), CO₂ (grade 5.0, Hoekloos), and PMMA (self-made, $M_{\rm w}$ and polydispersity 2.4×10^6 g/mol and 2.4, respectively). The MMA was distilled under vacuum to remove the hydroquinone inhibitor before use. To ensure that no reaction or recombination of polymer radicals could occur during the scission experiments, a radical scavenger (1,1-diphenyl-2-picrylhydrazyl, Aldrich) was added to the solution.

The composition of the liquid was calculated with the Lee-Kessler-Plöcker (LKP) equation of state (Table 1).24 In previous papers, 25,26 it was shown that the LKP-eos is the most

Table 1. Experimental Conditions for the Ultrasound-Induced Scission Experiments^a

T (°C)	P (bar abs)	$X_{{ m CO}_2}$	$I_{\mathrm{US}} (\mathrm{W/cm^2})$
20	1.0	0.02	31
20	4.0	0.08	59
20	7.0	0.14	70

^a The CO_2 fraction (X_{CO_2}) of the liquid was calculated with the Lee-Kessler-Plöcker equation of state.



High-pressure ultrasound reactor for scission Figure 2. reactions.

suitable equation of state to describe CO₂-monomer systems. The LKP model uses pure component parameters, such as the critical temperature, the critical pressure, and the acentric factor (ω) , and one binary interaction parameter. The interaction parameter for the MMA/CO₂ system appears to be 1.08 at 20 °C.²⁷ The influence of the polymer on the CO₂ fraction $(X_{\rm CO_2})$ in the liquid at a certain pressure was assumed to be negligible. The density of the liquid was approximately 0.945 kg/L at 20 °C.24

Scission Experiments. Ultrasound-induced scission was studied in a thermostated 175 mL high-pressure vessel at 20 °C (Figure 2). Sonification of the solution was performed using 20 kHz ultrasound, which was produced by a Sonics and Materials VC-750 generator. A ½-in. full-wave titanium probe was applied to couple the piezoelectric transducer to the liquid. To allow for an accurate comparison, the total configuration and the ultrasound amplitude (75 µm) were kept constant during all the experiments.

The experimental procedure of the ultrasound-induced scission experiments was as follows. The reactor was filled with 150 mL of bulk MMA or CO₂-expanded MMA, in which a known amount of PMMA was dissolved. To ensure the presence of nuclei for cavitation bubbles, CO2 was bubbled through a 3-mm tube into the reaction mixture with a flow rate of 2.0 mL/s during the experiments. This bubbling did not strip the MMA from the liquid phase because a condenser was placed on top of the reactor. As a result of the high solubility of CO2 in MMA, bubbling CO₂ through bulk MMA at atmospheric conditions already resulted in a CO₂ fraction of 0.02 (Table 1). In the CO₂-expanded MMA systems, a two-phase system was applied. The CO₂ gas phase was in equilibrium with the liquid CO₂/MMA phase. By adjusting the CO₂ pressure, the fraction of CO₂ could be controlled.

Analysis. During the scission reaction, samples were taken and analyzed by gel permeation chromatography (GPC), which was calibrated against polystyrene standards. The MWDs of the PMMA samples were calculated with Mark-Houwink parameters (a = 0.719 and $K = 9.44 \times 10^{-5} \text{ m}^3$).²⁸ The modeling procedure for the MWDs was as follows: the

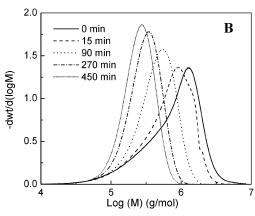


Figure 3. Experimental (A) and simulated (B) molecular weight distributions of an ultrasound-induced scission reaction of 4 wt% PMMA in bulk MMA ($X_{\text{CO}_2} = 0.02$).

Table 2. Experimental and Modeling Results of Ultrasound-Induced Scission in Bulk MMA ($X_{\text{CO}_2} = 0.02$) at Different Polymer Concentrations

wt% PMMA	$X_{{ m CO}_2}$	$M_{ m lim,experiment} \ (10^5 m g/mol)$	$M_{ m lim,model} \ (10^5 m g/mol)$	$A (10^{-5} 1/s)$	$\frac{B}{(10^{-11} 1/s)}$	$\frac{C}{(10^6)}$	$2\mathrm{e}^{-n/C}$ at 10^6 g/mol
0.1	0.02	1.0	0.9	-2 ± 1	2.4 + 0.1	0.2 + 0.1	1.9
2.0	0.02	1.9	0.9	-0.1 + 0.1	0.13 + 0.01	5 ± 1	2
4.0	0.02	1.9	1.0	-0.1 + 0.1	0.10 + 0.01	5 ± 1	2

experimentally obtained MWDs were fitted using PREDICI. The applied model in this fitting procedure is described in detail in a previous paper. ¹⁸ In this way, the molecular-weight-dependent scission rate constants $(k_d(n,m))$ could be determined on the basis of the constants A,B, and C (see eq 5). Subsequently, the $M_{\rm lim}$ was calculated and compared with the experimentally observed $M_{\rm lim}$.

The ultrasound intensities were determined at an amplitude of 75 μ m for the different conditions by calorimetry and are given in Table 1. These values give an indication of the ultrasound power transferred to the liquid in the high-pressure reactor. The heat flow (Q) that was generated by ultrasound in the calorimeter was calculated on the basis of the surface area of the reactor wall (A_r) , the overall heat transfer coefficient (U), and the difference between jacket (T_j) and reactor temperature (T_r) . Subsequently, the ultrasound intensity $(I_{\rm US})$, was calculated by dividing the heat flow by the surface area of the ultrasound horn $(A_{\rm US})$ according to eq 6.30

$$I_{\rm US} = \frac{Q}{A_{\rm US}} = \frac{UA_{\rm r}(T_{\rm r} - T_{\rm j})}{A_{\rm US}}$$
 (6)

Results and Discussion

In this study, the effect of the CO₂ fraction and the polymer concentration on the scission kinetics has been determined by comparing experimentally obtained MWDs with fitted MWDs using PREDICI. First, the ultrasound-induced polymer scission simulations will be discussed briefly. Next, the influence of the liquid viscosity on the scission kinetics has been determined in bulk MMA. To quantify the influence of CO2 on the implosion velocity independently from the CO2 antisolvent effect, first, the influence of the CO₂ fraction on the fracture rate has been measured at low polymer concentrations. Subsequently, the antisolvent effect of CO₂ on the scission kinetics has also been studied for more concentrated polymer solutions. Finally, the application of CO₂ to control the molecular weight during ultrasound-induced scission is discussed.

Simulations of Ultrasound-Induced Polymer Scission. Figure 3A shows the development of the MWD for a scission experiment in bulk with 4 wt% PMMA. A decrease in polydispersity and molecular

weight is obtained by irradiation of the solution with ultrasound. The modeled MWDs for this scission reaction are shown in Figure 3B. In a previous study, where monodisperse polymer samples have been used in bulk MMA, the width of the sigmoidal distribution function (eq 4) has been determined, which appears to be $\sigma(n) = 0.13n.^{18}$

In general, the model describes the scission reaction well, as only minor differences between the experimental and the model results can be observed. The calculated MWD for t=15 min shows a shoulder at the high-molecular-weight side. This is caused by the grid division in PREDICI because the calculation steps increase at higher molecular weight. Table 2 presents the limiting molecular weight and the constants and the exponent for the scission rate constant $k_{\rm MW}(n)$ as given in eq 5.

Influence of Viscosity on Transient Cavitation. The liquid viscosity has a major influence on the implosion velocity and, consequently, on the scission rate and $M_{\rm lim}$. The experimentally observed development of $M_{\rm w}$ at different polymer concentrations in bulk MMA is plotted in Figure 4A. In contrast to the solutions with 2.0 and 4.0 wt%, the low polymer concentration of 0.1 wt% is below entanglement and the coils of the polymer chains do not touch each other. Figure 4B shows the scission constants (k_d) , determined by the simulations as a function of molecular weight at different polymer concentrations in bulk MMA. It should be noted that, with decreasing molecular weights, the reaction time increases. A much lower scission rate constant is observed at higher polymer concentrations, which is a result of the higher viscosity. This decrease is the most distinct between polymer concentrations of 0.1 and 2 wt%, whereas the difference between 2 and 4 wt% dissolved polymer is small. Limiting molecular weights of 1.0×10^5 , 1.9×10^5 , and 1.9×10^5 g/mol are obtained for scission in bulk MMA for PMMA at concentrations of 0.1, 2, and 4 wt%, respectively (Table

With respect to simulations using PREDICI, Table 2 shows that the calculated $M_{\rm lim}$ value at low polymer

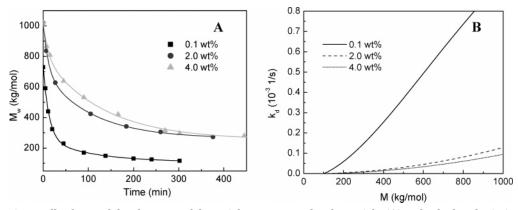


Figure 4. Experimentally observed development of the weight average molecular weight (A) and calculated scission rate constants (B) for ultrasound-induced scission at different PMMA concentrations in bulk MMA ($X_{CO_2} = 0.02$).

concentration is in good agreement with the experiment because the viscosity and, consequently, the cavitation intensity do not change during the experiment. Due to these constant experimental conditions, the scission model accurately predicts $M_{\rm lim}$ at low polymer concentrations. In contrast, at high polymer concentrations, the calculated $M_{\rm lim}$ values deviate from the experimentally obtained M_{lim} , approximately by a factor of 2. This is caused by the constant cavitation intensity in time used in the PREDICI model. This would imply that the viscosity of the solution remains constant during sonification. Experimentally, however, the viscosity changes significantly during the scission experiments of the more concentrated polymer solutions, as a result of which, the largest deviation in predicted $M_{\rm lim}$ is observed at the 2 and 4 wt% polymer solutions. In contrast to $M_{\rm lim}$, the scission model accurately predicts the initial scission rate constants at higher polymer concentrations.

In Table 2, the molecular-weight-dependent scission constants at different polymer concentrations are given. The exponent $(2e^{-n/C})$ given in Table 2 describes the dependence of the scission rate on the chain length (eq 5). The scission rate is expected to be a quadratic function of the length of the chain (eq 1). Although this is correct for low-molecular-weight polymers, from the simulations with 0.1 wt% polymer, it is concluded that for higher-molecular-weight polymers (10⁶ g/mol), the scission rate depends on the chain length to the power 1.9 (Table 2). This difference is significant, as described previously. 18 We assume that this deviance is probably a consequence of overlapping velocity profiles from neighboring cavitation bubbles (Figure 5). As a result, the volume in which a polymer with a certain molecular weight can break is reduced. The scission rate of a polymer chain is proportional to this volume, as mentioned in the theoretical section. In general, highmolecular-weight polymers can fracture further away from the bubble interface, and therefore, their scission rate will be more susceptible to this volume decrease. Consequently, the deviation from the theory (a quadratic function) is more distinct at higher molecular weights. Additionally, at higher cavitation intensities (lower viscosities), the radial velocity profile extends further in the liquid, which will cause more overlap with neighboring bubbles and, thus, a larger decrease in effective volume. As a result, a less-than-quadratic function is obtained at lower polymer concentrations.

Influence of CO₂ Fraction on Transient Cavitation. The influence of the CO₂ fraction on the implosion velocity has been measured independently from the

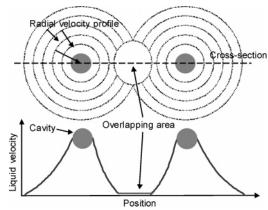


Figure 5. Schematic representation of overlapping velocity profiles around two imploding cavitation bubbles and the resulting radial velocity.

antisolvent effect by performing the experiments at low polymer concentrations (Table 3). At low polymer concentrations, the influence of the CO₂ fraction on the viscosity can be neglected because the polymer chains are not entangled. Figure 6A shows the development of $M_{\rm w}$ in time for scission experiments with a low polymer concentration at different CO2 fractions. With an increasing CO_2 fraction, the scission rate decreases (slope of Figure 6A) and the M_{lim} increases. The change in scission kinetics is a result of the lower strain rate produced upon collapse at higher CO₂ fractions. This is not an antisolvent effect because the CO₂ fraction has almost no influence on the viscosity in a 0.1 wt% polymer solution. However, the lower scission rate and higher M_{lim} at higher CO₂ pressures is a result of the higher vapor pressure in the cavitation bubble, which reduces the implosion velocity (cushioning effect) and, consequently, reduces the drag force acting on the polymer chain, see eq 1.

Figure 6B presents the overall scission constants at different CO2 fractions as a function of the molecular weight. From the simulations, it is concluded that for higher-molecular-weight polymers (10⁶ g/mol) the scission rate depends on the chain length to the power 1.9, 1.93, and 2 at 1, 4, and 7 bar, respectively (Table 3). This deviance from the theory is probably a consequence of overlapping velocity profiles from neighboring cavitation bubbles (Figure 5), as mentioned in the previous section. The larger deviation at lower pressures is caused by the higher implosion velocities. At lower CO₂ fractions, the radial velocity profile extends further in the liquid, which will cause more overlap with neigh-

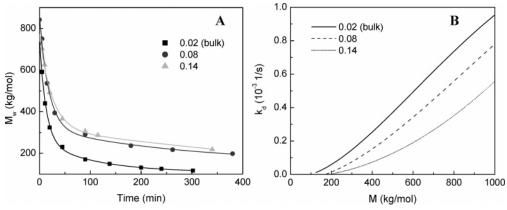


Figure 6. Experimentally observed development of the weight average molecular weight (A) and calculated scission rate constants (B) for ultrasound-induced scission at 0.1 wt% polymer and CO_2 fractions of 0.02, 0.08, and 0.14.

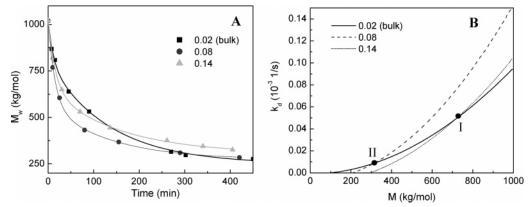


Figure 7. Experimentally observed development of the weight average molecular weight (A) and calculated scission rate constants (B) for ultrasound-induced scission at 4.0 wt% polymer and CO_2 fractions of 0.02, 0.08, and 0.14.

Table 3. Experimental and Modeling Results of Ultrasound-Induced Scission at Different CO₂ Fractions and 0.1 wt% Polymer

wt% PMMA	$X_{ m CO_2}$	$M_{ m lim,experiment} \ (10^5 m g/mol)$	$M_{ m lim,model} \ (10^5 m g/mol)$	$A (10^{-5} 1/s)$	$\frac{B}{(10^{-11}1/\mathrm{s})}$	$\frac{C}{(10^6)}$	$2\mathrm{e}^{-n/C}$ at $10^6\mathrm{g/mol}$
0.1	0.02	1.0	0.9	-2 ± 1	2.4 + 0.1	0.2 + 0.1	1.9
0.1	0.08	1.6	1.6	-4 ± 1	1.5 + 0.1	0.3 + 0.1	1.93
0.1	0.14	1.8	1.8	-2 ± 1	0.60 + 0.01	5 ± 1	2

Table 4. Experimental and Modeling Results of Ultrasound-Induced Scission at Different ${\rm CO_2}$ Fractions and 4.0 wt% Polymer

wt% PMMA	X_{CO_2}	$M_{ m lim,experiment} \ (10^5 m g/mol)$	$M_{ m lim,model} \ (10^5 m g/mol)$	$A (10^{-5} 1/s)$	$B \over (10^{-11} 1/s)$	$\frac{C}{(10^6)}$	$2\mathrm{e}^{-n/C}$ at 10^6 g/mol
4.0	0.02	1.9	1.0	-0.1 + 0.1	0.10 + 0.01	5 ± 1	2
4.0	0.08	2.6	2.1	-0.7 + 0.1	0.16 + 0.01	∞	2
4.0	0.14	3.0	2.9	-1 ± 1	0.12 + 0.01	∞	2

boring bubbles and, thus, results in a larger decrease in effective volume.

Influence of the Antisolvent Effect on the Scission Kinetics. Carbon dioxide reduces the viscosity of concentrated polymer solutions due to the antisolvent effect, 13 which induces a smaller gyration radius of the polymer or even polymer precipitation at higher CO_2 fractions. As a result, the scission rate constant is expected to remain virtually constant at higher polymer concentrations, whereas in bulk, the scission rate constant strongly decreases. Figure 7A shows the development of M_w in time for scission experiments for a 4 wt% polymer solution at several CO_2 fractions (Table 4). All the reaction mixtures are one-phase systems, as no precipitated polymer was observed in a high-pressure view-cell at these conditions. Figure 7B shows the scission rate constants for different CO_2 fractions, as

obtained from the model. It can be seen that the scission rate constant is initially (the right-hand side of point I and the right-hand side of point II) higher for the CO₂expanded liquids, in contrast to the bulk experiment (Figure 7B vs 6B). The negative influence of the CO₂ fraction on the implosion velocity at low polymer concentrations is thus counteracted at higher polymer concentrations by the CO2 antisolvent effect. As no precipitated polymer is observed, the higher scission rate is a consequence of the smaller gyration radius of the polymer chain, which results in a lower liquid viscosity. At the left-hand side of point I in Figure 7B, the scission rate in bulk is faster than the scission rate in CO_2 -expanded MMA at 7 bar ($X_{CO_2} = 0.14$). The difference with the right-hand side is caused by the larger viscosity decrease during the scission experiment in bulk MMA, which results in a larger increase in the

implosion velocity of the cavitation bubbles. This higher implosion velocity also results in a lower $M_{\rm lim}$ for the scission experiment in bulk MMA as compared to scission in CO₂-expanded MMA. The intersection point between CO_2 -expanded MMA at 4 bar ($X_{CO_2} = 0.08$) and bulk MMA (Figure 7B, II) occurs at a lower molecular weight because at 4 bar less CO₂ is dissolved in MMA, as compared to 7 bar. Therefore, the cushioning effect will be less pronounced at 4 bar CO₂.

As discussed in the previous sections, the CO₂ antisolvent effect and CO₂ cushioning effect have an opposite influence on the scission rate. Additionally, when substantial amounts of polymer are precipitated, it is possible that the breakage of polymer chains in solution is limited by the dissolution rate of the polymer from the precipitated phase. These effects may cause an optimum in scission rate at increasing static pressure.

Control of Molecular Weight by Ultrasound-Induced Polymer Scission. In general, ultrasound can be applied as a clean alternative method to control polymer properties by tuning the polydispersity and MWD because polymer scission occurs in situ without the addition of peroxides. Especially for polymers where the high-molecular-weight part of the distribution has a disadvantageous effect on the product properties, ultrasound has significant application potential. The possibility to perform these reactions at higher polymer concentration by the addition of an antisolvent makes the ultrasound technique more feasible for industry. In another study,31 a preliminary process design is presented for the production of 10 kg/h pure PMMA in CO₂expanded MMA including polymer scission. This process appears to be especially interesting for the production of specialty products for biomedical purposes.

Moreover, by combining ultrasound with CO₂, the scission rate and $M_{
m lim}$ can independently be altered. In comparison with bulk experiments, scission in CO2expanded liquids results in a lower viscosity and, consequently, a higher scission rate, as well as a lower $M_{\rm lim}$ in concentrated polymer solutions. Above a given CO₂ fraction, the polymer starts precipitating, resulting in a constant scission rate. The CO₂ fraction influences the amount of polymer that is dissolved, the M_{lim} , and the scission rate. Additionally, an increase in the ultrasound intensity will result in a lower $M_{\rm lim}$ and a higher scission rate, whereas the sonication area only influences the scission rate. In summary, the ultrasound intensity, the sonication area, and the CO₂ fraction can thus independently influence the limiting molecular weight and the scission rate. Therefore, ultrasoundinduced polymer scission in CO2-expanded liquids allows for an accurate control of the scission kinetics and, hence, the polymer properties.

Conclusions

In this work, the influence of the CO₂ antisolvent effect on the viscosity and the resulting reaction kinetics of ultrasound-induced polymer modification have been studied. For this purpose, ultrasound-induced polymer scission experiments have been performed in CO₂expanded MMA, as well as in bulk MMA, at different PMMA concentrations. Modeling the development of the experimental MWDs in PREDICI has revealed the scission kinetics at different polymer and CO₂ concentrations. At low polymer concentrations, the scission rate decreases at higher CO2 pressures due to the negative influence of CO₂ on the cavitation velocity. However, at higher polymer concentrations, this effect is counteracted by the viscosity reduction induced by the antisolvent effect of CO₂, which results in faster scission in CO2-expanded MMA, as compared to bulk MMA. When CO₂ acts as an antisolvent, it is possible to alter the MWD of polymers by ultrasound in concentrated polymer solutions.

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